Discharge coefficient of semi-circular labyrinth side weir in subcritical flow

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ABSTRACT

Side weirs are flow diversion or intake devices that are widely used in irrigation and drainage networks and urban sewage systems. Labyrinth weirs have a proven hydraulic advantage due to their increased discharge rate at the same head for a given design condition. A labyrinth weir is defined as a weir crest that is not straight in plan-form. The increased sill length provided by labyrinth weirs effectively reduces upstream head to the particular discharge. In this study, a comprehensive laboratory study was conducted in a physical model of a semi-circular labyrinth side weir and the model is evaluated for three heights and three radii. In this particular study, the hydraulic effects of this shape of side weir in increasing discharge capacity was investigated. The discharge coefficient of semi-circular labyrinth side weirs and the analysed to find the influence of the dimensionless parameters of weir height (w/b), weir length (r/b), nape height ($(y_1 - w)/r$), and upstream Froude number (Fr_1) on discharge coefficient (C_d), based on the experimental conditions and previous studies. It was found that the discharge coefficient of classic side weir positioned on a straight channel. Additionally, reliable equations for calculating the discharge coefficient (C_d) of semi-circular labyrinth side weirs are proposed. According to the proposed equation, C_d depends on dimensionless parameters, which are the ratios of r/b, r/w, $(y_1 - w)/r$, $(y_2 - w)/r$ and Fr_1 .

Keywords: side-weir, discharge coefficient, semi-circular labyrinth side weir, subcritical flow.

INTRODUCTION

A side weir is an overflow weir set into the side of a main channel. Side weirs have been extensively used in hydraulic and environmental engineering applications. They typically are used for water level control in canal systems, diverting excess water into relief channels during floods, as storm overflows from urban sewage systems, and as a head regulator of distributaries.

A review of literature on side weirs indicates that their discharge coefficient has been extensively studied, due to its importance. Rectangular sharp-crested side weirs have been investigated extensively, including work for example by Ackers (1957), Collings (1957), Frazer (1957), Subramanya and Awasthy (1972), El-Khashab and Smith (1976), Uyumaz and Muslu (1985). The analyses of Ranga Raju et al. (1979) and Ramamurthy and Carballada (1980) indicated that bed slope and friction highly affect side weir outflow. Borghei et al. (1999) discussed the application of spatially varied flow in sharp-crested rectangular side weirs. The hydraulic behaviour and discharge coefficient of side weirs for the different types of weirs, main channels and flow conditions have been studied by many researchers, for example, Nandesamoorthy and Thomson (1972), Yu-Tech (1972), Cheong (1991), and Singh et al. (1994). Ghodsian (2003) studied supercritical flow in rectangular side weirs. Cosar and Agaccioglu (2004) studied the discharge coefficient of a triangular side weir, for weirs on both straight and curved channels. Aghayari et al. (2009) experimentally investigated the effect of height, width and side weir crest slope

on the spatial discharge coefficient over broad-crested inclined side weirs under subcritical flow conditions in a rectangular channel. To study the discharge coefficient variation along a side weir, Swamee et al. (1994) used an elementary analysis approach to estimate the discharge coefficient in smooth side weirs through an elementary strip along the side weir.

Some of the proposed formulas for the estimation of discharge coefficient in rectangular side weirs are presented in Table 1.

Most of the presented equations for C_d depend on the Froude number of flow at the upstream section of the side weir. In addition, the majority of previous researchers have concentrated on investigating rectangular and triangular side weirs in straight channels. However, triangular and circular types of side weirs are also used in hydraulic and environmental engineering applications.

Labyrinth side weirs have different shapes, such as triangular, trapezoidal and semi-circular, in plan view. A labyrinth side weir provides a longer effective length for a given overall side weir opening. In fact, the effective length is the weir crest length that is denoted by *l*. The increased sill length provided by the labyrinth side weirs effectively reduces upstream head to the particular discharge. They can therefore be used to particular advantage where the width of a channel is restricted and a weir is required to pass a range of discharges with a limited variation in upstream water level. Emiroglu et al. (2010 a, b) studied the discharge coefficient of sharpcrested triangular labyrinth side weirs on a straight channel.

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Reference	C _d	Fr ₁	L/b	<i>w/y</i> ₁	<i>w</i> (m)
(1)	(2)	(3)	(4)	(5)	(6)
Subramanya & Awasthy (1972)	$C_{d} = 0.864 \left\{ \frac{1 - F_{1}^{2}}{2 + Fr_{1}^{2}} \right\}^{0.5}$	0.2-0.85	0.2–1.0	0.2-0.96	_
Yu-Tech (1972)	$C_d = 0.415 - 0.148 Fr_1$	-	-	0.2-0.5	-
Nandesamoorthy & Thomson (1972)	$C_{d} = 0.288 \left\{ \frac{2 - Fr_{1}^{2}}{1 + 2Fr_{1}^{2}} \right\}^{0.5}$	_	_	_	0.0-0.6
Ranga Raju et al. (1979)	$C_d = 0.54 - 0.4 Fr_1$	0.1-0.5	0.1-0.7	_	0.2-0.5
Hager (1987)	$\left C_{d} 0.488 \left\{ \frac{2 + Fr_{1}^{2}}{2 + 3Fr_{1}^{2}} \right\}^{0.5} \right $	0.0-0.87	3.33	_	0.0-0.2
Cheong (trapezoidal channel) (1991)	$C_d = 0.3 - 0.14 F r_1^2$	0.28-0.78	0.5–1.64	0.42-0.85	-
Singh et al. (1994)	$C_d = 0.33 - 0.18Fr_1 + 0.49\frac{W}{Y_1}$		0.25-0.5	0.42-0.85	_
Borghei et al. (1999)	$C_{d} = 0.7 - 0.48Fr_{1} - 0.3\frac{w}{y_{1}} + 0.06\frac{L}{b}$	0.1-0.9	0.67-2.33	-	0.01-0.2

 TABLE 1

 Side-weir discharge coefficient equations presented in the literature for straight channels. Fr, is upstream Froude number at the beginning

They proposed dimensionless parameters for triangular labyrinth side weir discharge coefficients on a straight channel as:

$$C_d = f(Fr_1, \frac{L}{b}, \frac{L}{l}, \frac{w}{\mathcal{Y}_1}, \theta, \psi)$$
(1)

where: Fr_1 , C_d , w, L, b and y_1 are defined earlier, l is overflow length of the side weir, θ is the included angle of the triangular labyrinth side weir. The water nape deviation or deflection angle ψ , is defined the deflection of the side-weir nape from the water surface toward the weir side and is given as follows (Subramanya and Awasthy, 1972):

$$\sin(\psi) = \sqrt{1 - \left(\frac{V_1}{V_s}\right)^2} \tag{2}$$

where: V_s is velocity of flow dQs over the brink. The water nape deviation ψ takes different values for each fluid particle and varies with the Froude number, which changes along the side weir due to spilling over the side weir. The deviation angle increases towards the side weir when the Froude number in the main channel decreases in the downstream direction. Some researchers have also mentioned that the dimensionless length of the side weir (L/b) includes the effect of the deviation angle on the discharge coefficient (Subramanya and Awasthy, 1972; El-Khashab, 1975; El-Khashab and Smith, 1976; Agaccioglu and Yüksel, 1998; Borghei et al. 1999). Therefore, the deviation angle ψ is not included in the side weir discharge coefficient equations in the literature. Emiroglu et al. (2010 a) reported that the discharge coefficient of the triangular labyrinth side weirs is 1.5–4.5 times higher than that of the rectangular side weir.

Unlike labyrinth spillways, one or two cycles are mostly sufficient for the labyrinth side weirs to deliver the required discharge. Therefore, there is a need to study different shapes of labyrinth side weirs in order to obtain comprehensive knowledge on this subject. Undoubtedly, an accurate and useful equation for semi-circular labyrinth side weirs with one cycle is required by hydraulic and environmental engineers. This study mainly aimed to evaluate various hydraulic characteristics of a semi-circular labyrinth side weir in one cycle by conducting experiments across a wide range of values for important parameters and to develop a useful and reliable equation for them in practical applications. This paper investigates the discharge coefficient of semi-circular labyrinth side weirs under subcritical flow condition and, in particular, the effect of the upstream Froude number Fr_1 and dimensionless parameters (w/b), (r/b), (w/y_1) on the discharge coefficient, given the experimental conditions and results of previous studies. The discharge coefficient of semi-circular labyrinth side weirs was determined based on dimensionless parameters, which are the ratios of r/b, r/w, b/y_2 , $(y_1-w)/r$, $(y_2-w)/r$ and Fr_1 . Also, the longitudinal velocity of semi-circular labyrinth side weirs has been studied for the subcritical flow in detail.

DISCHARGE COEFFICIENT FUNCTION

The flow over a side weir is categorized as a spatially varied flow condition. The present study dealt mainly with the application of the energy principle in the analysis of flow over side weirs. The concept of constant specific energy (De Marchi, 1934) is quite often adopted in studying the flow characteristics of side weirs (Emiroglu et al., 2010a, b; Singh et al., 1994; Hager, 1987; Subramanya and Awasthy, 1972). De Marchi (1934) was one of the first researchers to provide equations for flow over side weirs. He assumed that the discharge over the side weir per unit length, q, is:

$$q = \frac{dQ}{ds} = C_d \sqrt{2g} (y - w)^{1.5}$$
(3)

where: Q = discharge in the main channel at section *s*; *s* = distance from the beginning of the side weir; *q* (or = dQ/ds) = discharge spilling for per unit length of the side weir; *g* = acceleration due to gravity; *w* = height of the side weir;

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a) Longitudinal cross-section for the rectangular side weir



Figure 1 Definition sketch of subcritical flow over a rectangular side weir

y = depth of flow at section *s*; and $C_d =$ discharge coefficient of the rectangular side weir (Fig. 1). Thus the side weir discharge equation can be written as:

$$Q_w = 2/3C_d L \sqrt{2g} (y - w)^{1.5}$$
(4)

where: Q_w is total flow over side weir, C_d is the discharge coefficient and is dimensionless, L is the width of side-weir, and y and w are main channel depth and height of weir crest, respectively.

Understanding the physical relationship between discharge coefficient and independent non-dimensional variables is very important to be able to construct an appropriate function. In this study, a dimensionless analysis based on Buckingham's theorem was used to find non-dimensional variables. After the determination of dimensionless variables and elimination of some of the parameters, the functional discharge coefficient can be obtained by:

$$C_{d} = \phi \left(\frac{w}{y_{1}}, \frac{b}{y_{1}}, \frac{V_{2}}{y_{1}}, Fr_{1}, \frac{(y_{1} - w)}{r}, \frac{r}{b}, \frac{w}{y_{2}}, \frac{b}{y_{2}}, \frac{(y_{2} - w)}{r}, \frac{r}{w} \right)$$
(5)

where: Fr_1 is approach Froude number upstream of the side weir, w is side weir height, l is the effective side weir length, r is the radius of semi-circular labyrinth side weir, b is the main channel width, V_1 and V_2 are the mean velocities in the upstream and downstream sections of the weir crest, respectively, y_1 and y_2 are the depth of flow at the upstream and downstream end of the side weir in the centreline of the main channel. A longitudinal section and plan of a semi-circular labyrinth side weir with different variables are presented in Fig. 2.

EXPERIMENTAL SET-UP

Semi-circular labyrinth side weir experiments were conducted at the hydraulic laboratory of Shiraz University, Shiraz, Iran. The experimental set-up consisted of a main channel and the discharge collection channel. The main channel was 15 m long with a rectangular cross-section. The main channel had

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Figure 2

Definition sketch of subcritical flow over a semi-circular labyrinth side weir

internal width b = 0.50 m and 0.40 m depth with the bed slope at 1/1 000. The length of the channel from the inlet section to the beginning of the side weir was 7 m. A sluice gate is fitted at the end of the main channel to control flow depth. The lateral channel was parallel to the main channel and was 8 m long, 0.60 m wide and 0.40 m deep with a horizontal invert. A schematic representation of the semi-circular labyrinth side weir set-up is shown in Fig. 3. The main and lateral channels were made of brick masonry and plastered with cement. Two pre-calibrated standard V-notches were used to measure the discharge of the downstream end of the lateral and main channels. The side weirs were made of mild steel plate, the top edge being suitably bevelled to get a sharp crest. The side weir was installed flush with the main channel wall. To measure the water depth, a point gauge with an accuracy of ± 0.1 mm was used. Water depth was measured in the centreline axis of the channel length. Water depth measurements were made using a handmade type of frame, which can move in both xand y directions on a rail. Velocity was measured by using an electromagnetic velocity meter with an accuracy of 1 mm/s (model LP1100). To understand the hydraulic behaviour of the semi-circular labyrinth side weir, two sets of velocity measurements are taken at the side weir sections of the flume. The first set was taken with low overflow condition $(Q_{u} = 0.2 Q_{1})$. The second was taken with the high overflow condition $(Q_w =$ $0.65 Q_1$). Longitudinal velocities are measured at the point showed in Fig. 4. Experiments were conducted for open-end and semi-closed-end conditions. These experiments were conducted for subcritical flow, stable and free overflow conditions. Coleman and Smith (1923) pointed out that minimum nape height should be less than 19 mm because of the surface tension over the weir crest. Therefore, minimum nape height was taken as 20 mm. The experiments were conducted for 3 heights and 3 radii of semi-circular labyrinth side weir (w = 0.05, 0.10, 0.15 m and r = 0.18, 0.29, 0.40 m). In case of semi-circular labyrinth side weir with radius and height of 0.29 m and 0.10 m, respectively, the equivalent classical rectangular side weir was also tested and measured discharge coefficients were compared. The range of parameters is given in Table 2.

TABLE 2								
Range of variables tested								
<i>r</i> (m)	<i>w</i> (m)	<i>b</i> (m)	r/b	<i>w/y</i> ₁	Q ₁ (L/s)	Fr ₁		
0.18, 0.29, 0.40	0.05, 0.10, 0.15	0.50	0.36, 0.58, 0.80	0.24-0.89	2.9-47.0	0.06-1.00		



Figure 3 Schematic view of semi-circular labyrinth side weir located on straight channel

EXPERIMENTAL RESULTS AND ANALYSIS

Vertical profiles of longitudinal velocities in the main channel

Although this study did not anticipate investigating the velocity distribution in detail, an important feature was to observe the shape of the curve for maximum longitudinal velocity in order to know the effect of overflow on the main channel flow. Velocity measurements were conducted for both high and low overflow conditions ($Q_w = 0.65Q_1$ and $Q_w = 0.2 Q_1$) for the semi-circular labyrinth side weir with radius of r = 0.29 m (weir length of L = 0.58 m), weir height of w = 0.10 m and, discharge of main channel of $Q_1 = 30$ L/s.

Vertical velocity profiles at the point of C_1 , C_2 , C_3 , C_4 and C_5 ; D_1 , D_2 , D_3 , D_4 and D_5 ; E_1 , E_2 , E_3 , E_4 and E_5 were measured. As can be seen in Fig. 4, velocity profiles at the C–C, D–D and E–E sections of the straight channel (v/v_m) are plotted versus z/z_0 . The velocities are normalized by dividing them with the maximum longitudinal velocity, which is measured along the channel, and the depth is normalized by dividing the measured height with the depth of flow in the same run of velocity measurements, in which v_m is maximum velocity, and z and z_0 are depths where z_0 is equal to y (depth of flow).

Typical vertical profiles of longitudinal velocities along the semi-circular labyrinth side weir are shown in Fig. 5. Comparing the results of this study with those of Emiroglu et al. (2010 a), the same findings can be observed. The v/v_m ratio decreases towards the inner bank due to lateral flow; the v/v_m ratio for the inner bank is lower than that of outer bank. The inner bank has been delineated by Longitudinal Profile 5 of the flow through the main channel and the outer bank has been delineated by Longitudinal Profile 1, where the side weir is set





Figure 4 Places that longitudinal velocities were measured

in Fig. 4. In section E-E, velocity in the main channel decreases due to lateral flow at the downstream end of the side weir and thus flow depth at the downstream end become higher than that of the upstream end. The velocity value tends towards zero at the downstream end of the side weir. The effect of lateral flow is significant, especially for high overflow conditions.

C₄ coefficient

Experiments were conducted to determine the discharge coefficient of semi-circular labyrinth side weirs. The discharge coefficient was computed using the De-Marchi equation (Eq. 4):

$$L = \frac{3b}{2C_d} (\phi_2 - \phi_1) \tag{6}$$

where:

$$\phi \frac{2E_s - 3w}{E_s - w} \left(\frac{E_s - y}{y - w}\right)^{0.5} - 3\sin^{-1} \left(\frac{E_s - y}{E_s - w}\right)^{0.5} \tag{7}$$

and: *b* is the main channel width, *L* is the weir crest width, *y* is the depth of flow, *w* is side weir height and E_s is the specific energy that is assumed to be constant along the length of the side weir.

The crest length of the semi-circular labyrinth side weir is always longer than that of the classical rectangular side weir. In other words, if the crest length of the rectangular side weir is considered as *L* and the crest length of the semi-circular labyrinth side weir with one cycle is considered as $l = \pi \cdot r = \pi \cdot L/2$, then l/L = 1.57 can be calculated for the semi-circular labyrinth side weir. As we are obliged to use the length (width) of side weir (*L*) in Eq. 6, whereas the real length crest of a semicircular labyrinth side weir is *l* and is 1.57 times longer than width of side weir (*L*), then it is expected to obtain a higher discharge coefficient (C_q) in comparison with classic side weirs.



Figure 5 Typical vertical profiles of longitudinal velocities along the side weir



Figure 6 Comparison of the C_d values for a semi-circular labyrinth side weir with the equivalent classical rectangular side weir plotted against Fr_{r_1} (for w/b = 0.2 and L/b = 1.16)

Since the weir length is excessive, more intensive secondary flow occurs. Thus, intensive secondary flow can also result with higher flow rate (Kaya et al., 2011). Therefore, the main purpose of a labyrinth side weir is to create a higher crest length as well as other hydraulic effects such as more severe secondary flow (secondary considerations).

Dimensional analysis allows description of all important factors involved in the phenomenon expressed in dimensionless components. A multicollinearity test was used to analyse the data (variable). Results showed that for a sharp-crested semi-circular labyrinth side weir, the discharge coefficient is a function of:

$$C_{d} = f\left(\frac{r}{b}, Fr_{1}, \frac{r}{w}, \frac{b}{y_{2}}, \frac{(y_{1} - w)}{r}, \frac{(y_{2} - w)}{r}\right)$$
(8)

or

$$C_{d} = f\left(\frac{r}{b}, Fr_{1}, \frac{r}{w}, \frac{(y_{1} - w)}{r}, \frac{(y_{2} - w)}{r}\right)$$
(9)

with regard to linear or non-linear regressions, respectively.

The discharge coefficient values for a semi-circular labyrinth side weir with r = 0.29 m and w = 0.10 m are compared with those of the equivalent classical rectangular side weir in Fig. 6. The values of C_d were plotted against Fr_1 . It can be seen in Fig. 6 that the values of C_d for the semi-circular labyrinth side weir with L/b = 1.16 and w/b = 0.2, are significantly higher than those of the rectangular side weir that was measured in this study. The primary reason for the increase in the discharge capacity of the semi-circular labyrinth side weir can be attributed to the increase in the overflow length of the side weir.

The discharge coefficient values for the semi-circular labyrinth side weir (r/b = 0.36, 0.58, 0.80) are compared with those of Ranga Raju et al. (1979), Subramanya and Awasthy (1972) and Nandesamoorthy (1972) for rectangular side weirs. As can be seen in Fig. 7, the values of C_d for the semi-circular labyrinth side weirs are significantly higher than those suggested in previous studies. Based on the literature and dimensional analysis results, semi-circular labyrinth side weirs were tested for three different heights and radii, to investigate the effect of different heights and radii on discharge coefficient. To study the effect of parameter upstream Froude Number (Fr_1) and dimensionless parameter of the nape height ($(y_1-w)/r$) on the discharge coefficient, the values of C_d are plotted against Fr_1 and $(y_1-w)/r$ in Figs 8 and 9, respectively, for all of the experimental data.

As can be seen in Fig. 8, the discharge coefficient decreases when the Froude number increases. This decreasing trend is counter to that reported by Emiroglu et al. (2010 a) for triangular labyrinth side weirs, but is similar to the recent study about semi-elliptical labyrinth side weirs reported by Kaya et al. (2011). Scattering of the data is attributed the influence of the effective parameters, such as r/b. Fig. 9 reveals that C_1 almost has an indistinct trend versus the dimensionless parameter of the nape height $(y_1 - w)/r$. Although the scattering of data is very apparent, it shows a moderate descending trend in C, values as $(y_1 - w)/r$ increases. Figure 10 shows the effect of weir height on discharge coefficient of a semi-circular labyrinth side weir in different ratios of w/b and the ratio of r/b = 0.8. As can be seen in Fig. 10, C_d takes higher values when the dimensionless parameter of the weir's height *w/b* increases. Kaya et al. (2011) pointed out that the effect of exerting different heights (in the form of dimensionless parameters w/b or w/y_1) on C_d can be explained by the existence of a discontinuity region.



Fiaure 7

Comparison of the C_d values of semi-circular labyrinth side weir with the equations of Ranga Raju et al. (1979), Subramanya and Awasthy (1972), and Nandesamoorthy (1972) for three ratios of r/b = 0.8, 0.5, 0.36 (a, b, c)

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Figure 10 C_{a} versus Fr, for different ratios of w/b = 0.1, 0.2, 0.3 in r/b = 0.8

This discontinuity region has a strong secondary motion next to the boundary of the weir-side. The intensity of this secondary motion next to the boundary depends on the crest height of the side weir and decreases with an increase in height, due to the friction of the weir surface.

The variation of discharge coefficient of the semi-circular labyrinth side weir is presented for a ratio of w/b = 0.3 for r/b = 0.36 and r/b = 0.8 (dimensionless parameter of weir's length) in Fig. 11. It is clearly observed from Fig. 11 that discharge coefficients of the semi-circular labyrinth side weirs have higher values at higher ratios of r/b. In other words, higher C_d values are obtained at high r/b ratios due to an increase in the intensity of secondary flow created by lateral flow. El-Khashab and Smith (1976) also pointed out that the secondary flow condition due to lateral flow is dominant when a side weir is relatively long (i.e. L/b > 1). A secondary flow is a relatively minor flow superimposed on the main flow which can be explained by the existence of a discontinuity region. The intensity of this secondary motion next to the boundary of the weir side depends on the crest height of the side weir and decreases with the crest height increase of the side weir due to the friction of the weir surface (Kaya et al., 2011).



Figure 9 C_d versus different $(y_1 - w)/r$ values



Figure 11 C_{a} versus Fr, for different ratios of r/b = 0.36 and 0.8 in w/b = 0.3

Empirical correlations to predict discharge coefficient C_{d} were developed for semi-circular labyrinth side weirs, according to the result of dimensionless analysis that contains linear and non-linear equations. The resulting correlations are given in Eqs. 10 and 11.

$$\begin{split} C_{d} &= 0.7874 + 0.379 \left(\frac{r}{b}\right) - 0.3383 Fr_{1} - 0.0466 \left(\frac{r}{w}\right) + \\ 0.0276 \left(\frac{b}{y_{2}}\right) - 2.052 \left(\frac{y_{1} - w}{r}\right) + 0.161 \left(\frac{y_{2} - w}{r}\right) \end{split} \tag{10}$$

$$C_{d} &= 0.99 \left\{ (1.02 + 254.31 \left(\frac{r}{b}\right)^{37.5} - 0.5 Fr_{1}^{1.9} - 0.001 \left(\frac{r}{w}\right)^{2.38} - 2.915 \left(\frac{y_{1} - w}{r}\right)^{0.557} + 2.314 \left(\frac{y_{2} - w}{r}\right)^{0.533} \right\}^{1.023} \tag{11}$$

It should be noted that the discharge coefficient of the semicircular labyrinth side weir is calculated for a single cycle. To compare the computed C_d values with the measured values, using Eqs 10, and 11, Figs 12 and 13 are presented, respectively. From Figs 12 and 13, it is evident that the majority of the data points fall in the error band of 10%. Therefore, very good agreements between the measured values and those computed from the predictive equations are obtained.

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Figure 12 Comparison of measured C_d values with those calculated from Eq. 9

To evaluate the accuracy of the estimates produced using the obtained equations, the relative error (RE), root mean square error (RMSE) and index of agreement (d) criteria were used. These accuracy criteria are defined as:

$$RE = \frac{100}{N} \sum_{i=1}^{N} \left| \frac{X_{O_i} - X_{c_i}}{X_{O_i}} \right|$$
(12)

$$RMSE = \sqrt{\frac{1}{N} \sum_{i=1}^{N} (X_{O_i} - X_{c_i})^2}$$
(13)

$$d = 1 - \left[\frac{\sum_{i=1}^{N} (X_{O_i} - X_{c_i})^2}{\sum_{i=1}^{N} (\left| \overline{X}_{O_i} - X_{c_i} \right| + \left| \overline{X}_{O_i} - \overline{X}_{O_i} \right|)^2} \right]$$
(14)

where: N is the number of data points and X_{a} and X_{c} are the observed and computed values of discharge coefficient, respectively. Also \overline{X} is the mean of the observed values of the discharge coefficient. The computed values of RE, RMSE and d for the obtained C_d equations are presented in Table 3.

TABLE 3 Statistical criteria for presented equations					
Equation	%RE	RMSE	D		
0.94	0.067	7.14	9		
0.95	0.062	6.89	10		

CONCLUSIONS

Laboratory experiments of the semi-circular labyrinth side weir located on a straight channel were carried out to investigate the effect of different dimensionless parameters on the discharge coefficient. Empirical equations predicting the discharge coefficient of the semi-circular labyrinth side weirs were developed. According to the suggested equation (non-linear equation), C_{d} depends on dimensionless parameters that are r/b, Fr_1 , r/w, $(y_1 - w)/r$ and $(y_2 - w)/r$. The following conclusions can be drawn based on these findings:

The longitudinal velocity of water moving in active overflow condition from the inner bank towards the outer bank increases due to lateral flow



Fiaure 13 Comparison of measured C_d values with those calculated from Eq. 10

- The discharge coefficient of the semi-circular labyrinth side weir is higher than that of the rectangular side weir
- The discharge coefficient increases with increasing r/b ratio
- The discharge coefficient increases with increasing w/b ratio
- The proposed equations for C_{d} , the De-Marchi Coefficient, were shown to be reliable for the subcritical flow condition in this study. Although a non-linear equation is suggested to predict the values of C_d , the linear equation generates similar predictions for C_d .

NOTATION

- b width of channel
- С, side-weir discharge coefficient
- Fr_1 Froude number at upstream end of side weir
- y main channel depth
- flow depth at upstream end of side weir at channel center y_1
- flow depth at downstream end of side weir at y_2
 - channel center
- L length (width) of side weir
- weir crest length l
- height of weir crest w
- radius of semi-circular labyrinth side weir r
- Q discharge in the main channel
- Q_1 discharge at the upstream end of the side weir in the main channel
- Q, discharge at the downstream end of the side weir in the main channel
- Q_{w} total discharge overflowed the side weir
- E specific energy
- s distance alongside weir measured from upstream end of side weir (m)
- Vmean velocity in any section of channel
- $V_{\rm m}$ $V_{\rm r}$ maximum velocity of spiral flow along channel
- mean velocity of flow at upstream end of side weir
- V/V_{m} dimensionless velocity at any point
- local vertical coordinate z
- is equal to y Z_0
- z/z_{o} dimensionless depth at any point
- ψ deviation angle of flow

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